

Transition to turbulence in channel flow with spanwise magnetic field

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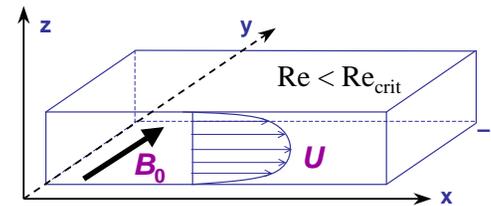
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MHD flow with spanwise magnetic field

Channel geometry



$$Re = \frac{aU}{\nu}$$

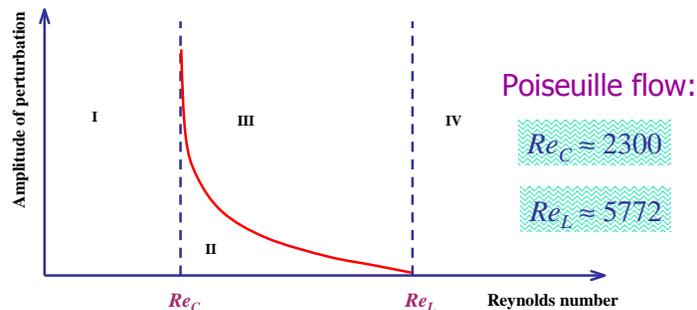
$$Ha = B_0 a \sqrt{\frac{\sigma}{\rho \nu}}$$

$$N = \frac{Ha^2}{Re} = \frac{B_0^2 a \sigma}{\rho U}$$

Question:

How does the transition to turbulence occur in the presence of a spanwise magnetic field?

Stability of shear flows



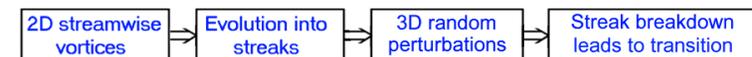
- Globally stable (initial growth of perturbations is possible)
- Conditionally stable: initial growth, then asymptotic decay
- Conditionally unstable: instability will occur
- Linearly unstable

Transition scenario (non-MHD case)

- Basic flow + 2D streamwise vortices (optimal disturbance: strong growth for finite time)
- At time t_{opt} of maximum growth: impose small 3D random perturbation

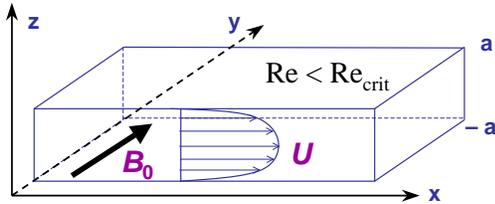
$$\mathbf{V} = U(z)\mathbf{e}_x + \mathbf{V}_2(y, z, t) + \epsilon \mathbf{V}_3(x, y, z, t)$$

- Route to transition:



(*streak breakdown*, Henningson, Reddy, Schmid)

Governing equations



$$Re = \frac{aU}{\nu}$$

$$Ha = B_0 a \sqrt{\frac{\sigma}{\rho\nu}}$$

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = -\nabla p + \frac{1}{Re} \Delta \mathbf{v} + \frac{Ha^2}{Re} (\mathbf{j} \times \mathbf{e}_0)$$

$$\mathbf{j} = -\nabla \Phi + \mathbf{v} \times \mathbf{e}_0, \nabla \cdot \mathbf{j} = 0, \Delta \Phi = \nabla \cdot (\mathbf{v} \times \mathbf{e}_0)$$

In linearized problem solutions are Fourier modes with respect to homogeneous directions x and y -> streamwise and spanwise wavenumbers α and β

Linear problem formulation

Linear evolution equations formulated for the following variables:
 η - vertical vorticity, w - vertical velocity, Φ - electric potential

- $\partial_t \eta + i\alpha U \eta - (D^2 - k^2) \eta / Re = i\beta w DU - i\beta N [D\Phi - (i\alpha Dw + i\beta \eta) / k^2]$
- $[\partial_t + i\alpha U - (D^2 - k^2) / Re] (D^2 - k^2) w = i\alpha w D^2 U + \beta^2 N w$
- $(D^2 - k^2) \Phi = (i\alpha D^2 w + i\beta D \eta) / k^2 - i\alpha w$
 $k^2 = \alpha^2 + \beta^2, D \equiv \partial / \partial z, N = Ha^2 / Re$

perturbation energy

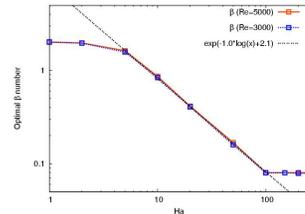
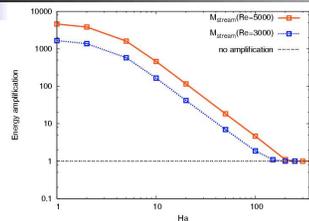
$$E(t) = \int_{-1}^1 (|Dw|^2 + k^2 |w|^2) dz + \beta^2 Re^2 \int_{-1}^1 |\eta|^2 dz$$

Energy amplification

$$G(t, \alpha, \beta, Ha, Re) = \max |E(t)| / |E(0)|$$

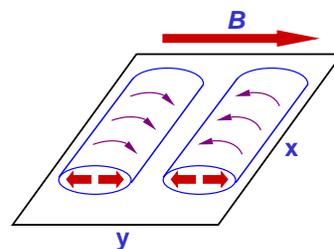
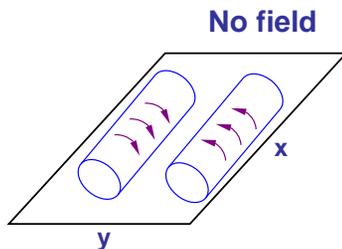
$$M(t, Ha, Re) = \sup_{\alpha, \beta} G(t, \alpha, \beta, Ha, Re)$$

Optimal streamwise vortices: damping of transient growth by magnetic field



Energy amplification M vs. Ha for Reynolds numbers $Re = 3000$ and 5000

Optimal spanwise wavenumber β vs. Ha for Reynolds numbers $Re = 3000$ and 5000



Arguments for the scaling of streamwise vortices

Linear results yield: $M_{stream} \approx Ha^2$ and $\beta_{stream} \approx Ha^{-1}$ for $5 < Ha < 100$

New scaled variables: $t' = t / Re, \eta' = \eta / (\beta Re), w' = w, \Phi' = \Phi / Re$

We find that

$$\partial_t D \eta' - (D^2 - \beta^2) D \eta' = iD (w' DU) - i\beta^2 Ha^2 \Phi'$$

Squire equation is responsible for the energy amplification - as in ordinary hydrodynamics
 $\Rightarrow M_{stream}(\beta, Ha, Re) = \beta^2_{stream} Re^2 F(R), R = \beta^2_{stream} Ha^2$ and $F(R)$ is the amplification given by:

$$\partial_t D \eta' - D^2 D \eta' = iD (w' DU) - iR \Phi'$$

$$[\partial_t - D^2] D^2 w' = R w$$

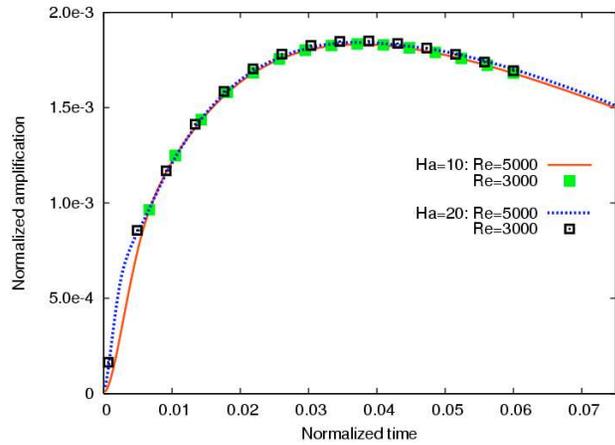
$$D^2 \Phi' = iD \eta'$$

Let $F(R)$ have a maximum at $R = R_c$ and decrease as $1/R^p$ with $p > 1$ for $R \rightarrow \infty$

\Rightarrow Three possible cases

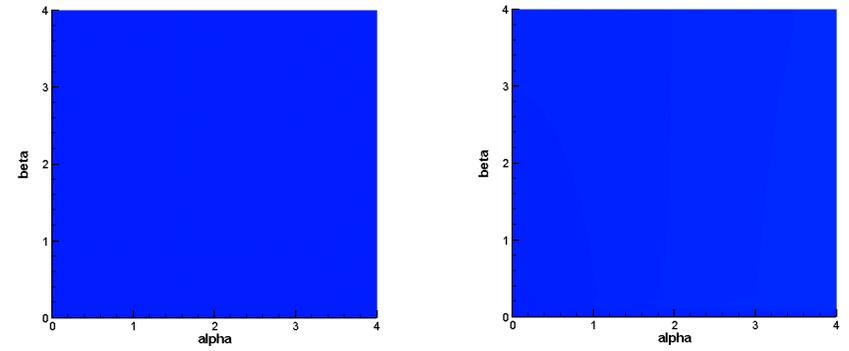
- If $\beta^2_{stream} \approx 1/Ha^n$ with $n > 1$ then: $R \rightarrow 0$ and $M_{stream} \rightarrow Re^2 / Ha^{2n} F(0)$
- If $\beta^2_{stream} \approx 1/Ha^n$ with $0 < n < 1$ then: $R \rightarrow \infty$ and $M_{stream} = Re^2 \beta^2_{stream} F(R \rightarrow \infty) \ll Re^2 / Ha^2$
 \Rightarrow decreases faster than $1/Ha^2$
- If $\beta^2_{stream} \approx 1/Ha$ then: $R = const$ and $M_{stream} = R_c F(R_c) Re^2 / Ha^2$
 \Rightarrow hypothesis 3 is correct since it results in the observed scaling

Streamwise vortices: damping and scaling



Energy amplification M vs. time T for $Re = 3000$ & 5000 and $Ha = 10$ & 20 : in rescaled units nearly collapse into one curve

General case: perturbations of arbitrary orientation



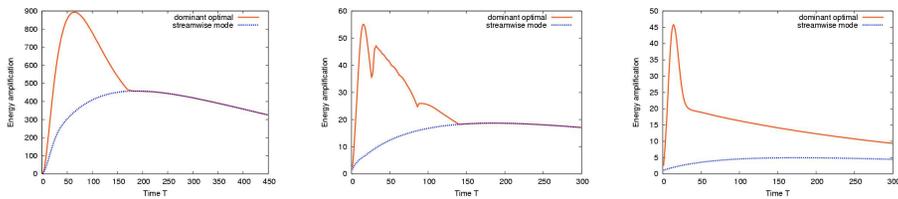
$Ha = 10$ & $Re = 5000$

$Ha = 50$ & $Re = 5000$

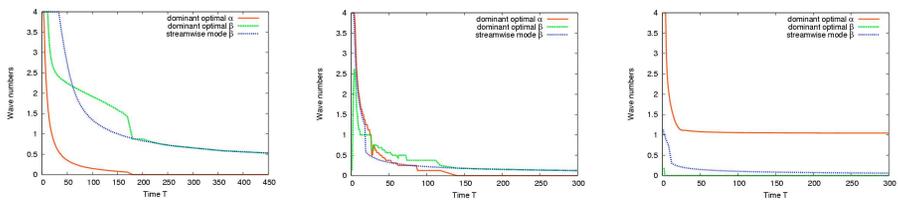
Isolevels of energy amplification $M(\alpha, \beta, T)$ in the entire (α, β) domain

Change of energy amplification and optimal wavenumbers vs. time T and Ha at $Re = 5000$

Energy amplification M for global and streamwise optimal modes



Optimal wavenumbers (α, β) for global and streamwise modes

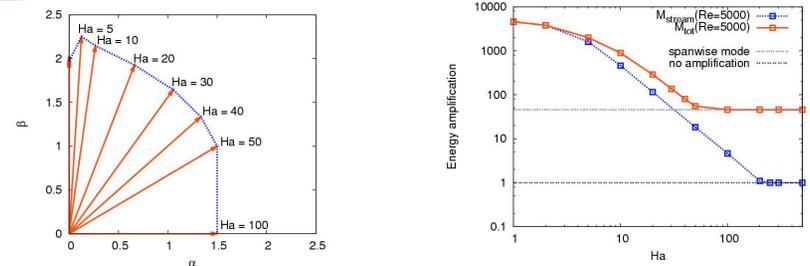


$Ha = 10$

$Ha = 50$

$Ha = 100$

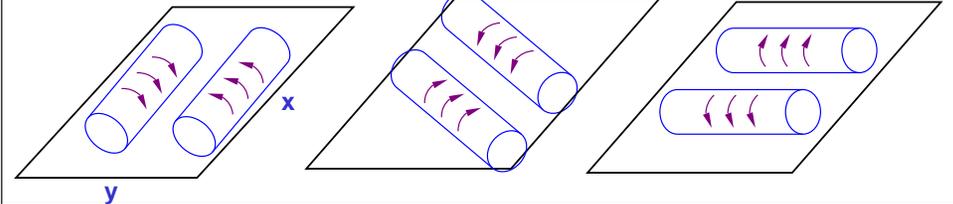
Optimal oblique modes at $Re = 5000$: summary



No field

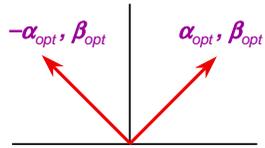
$Ha \approx 30 \dots 50$

$Ha \geq 100$

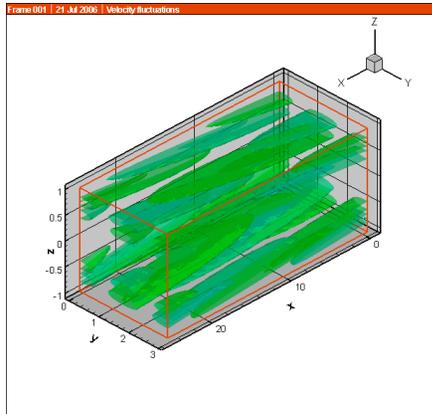
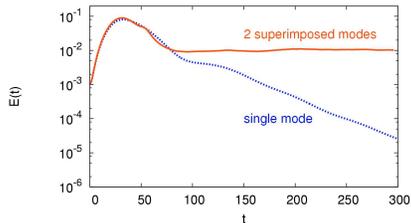


Superposition of optimal oblique modes as transition scenario

Optimal (α, β) mode and its symmetric $(-\alpha, \beta)$ counterpart provide the same linear growth



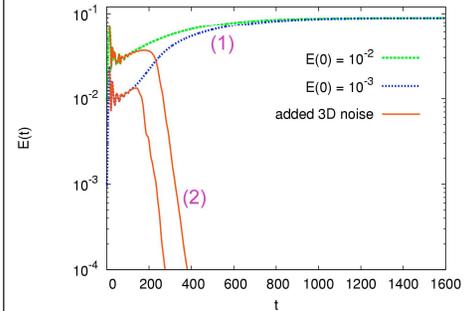
Non-linear interaction of superimposed oblique modes as a source of instability



Transition to turbulence by non-linear interaction of two superimposed oblique modes. $Re = 5000$ and $Ha = 10$, iso-surfaces of streamwise velocity perturbations

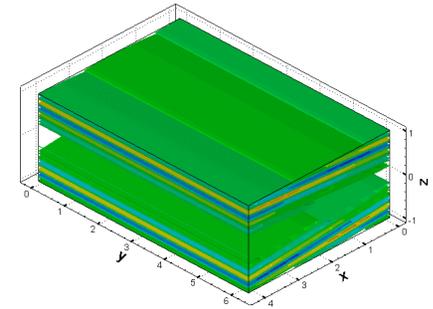
Transition at high Hartmann numbers ($Ha \geq 100$)

At high Ha purely 2D Orr-modes ($\alpha \neq 0, \beta = 0$) become optimal: no interaction with spanwise magnetic field



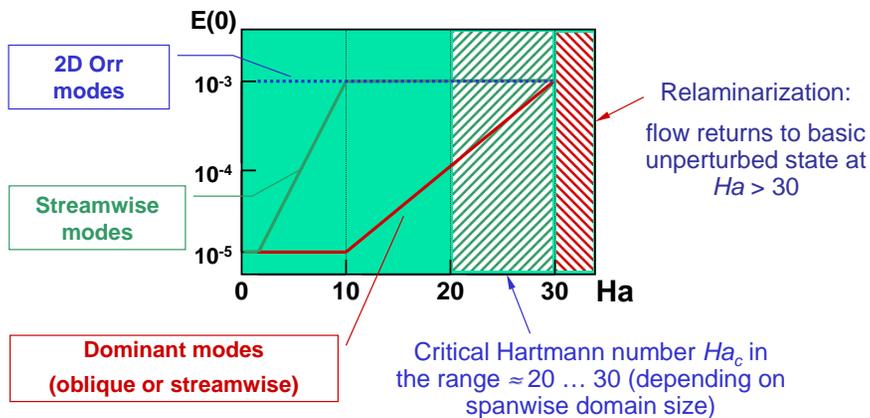
Evolution of Orr-modes at $Re = 5000$ and $Ha = 100$

- (1) If no noise added: flow remains purely 2D
- (2) 3D noise: triggers re-laminarization – Joule dissipation drains energy from 2D modes



Flow re-laminarization and return to unperturbed state triggered by 3D noise. $Re = 5000$ and $Ha = 100$, iso-surfaces of streamwise velocity perturbations

Stability thresholds for initial energy of optimal perturbations $E(0)$ versus Ha



Summary and outlook

Summary:

- streamwise modes strongly damped by magnetic field
- highest energy amplification for larger Ha for oblique modes
- oblique modes align with B as Ha increases
- superposition of symmetric oblique modes favorable for transition
- relaminarization of 2d time-dependent flow at high Ha

Things to do:

- classification of oblique modes
- transition studies with high aspect ratios
- comparison with experiments?

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