

# SETTLING VELOCITY OF INERTIAL PARTICLES

Marco Martins Afonso

Department of Physics of Complex Systems,  
Weizmann Institute of Science, Rehovot (Israel)

Antonio Celani (Nice) – Andrea Mazzino (Genova)

ETC11 Porto, 25\6\2007

# INERTIAL PARTICLES: MODEL AND LIMITATIONS

# INERTIAL PARTICLES: MODEL AND LIMITATIONS

## Model:

- ▶ given incompressible surrounding flow

# INERTIAL PARTICLES: MODEL AND LIMITATIONS

## Model:

- ▶ given incompressible surrounding flow
- ▶ single spherical particle

# INERTIAL PARTICLES: MODEL AND LIMITATIONS

## Model:

- ▶ given incompressible surrounding flow
- ▶ single spherical particle
- ▶ size smaller than (smallest active) scale of flow

# INERTIAL PARTICLES: MODEL AND LIMITATIONS

## Model:

- ▶ given incompressible surrounding flow
- ▶ single spherical particle
- ▶ size smaller than (smallest active) scale of flow
- ▶ low relative Reynolds number

# INERTIAL PARTICLES: MODEL AND LIMITATIONS

## Model:

- ▶ given incompressible surrounding flow
- ▶ single spherical particle
- ▶ size smaller than (smallest active) scale of flow
- ▶ low relative Reynolds number
- ▶ Stokes' viscous drag

# INERTIAL PARTICLES: MODEL AND LIMITATIONS

## Model:

- ▶ given incompressible surrounding flow
- ▶ single spherical particle
- ▶ size smaller than (smallest active) scale of flow
- ▶ low relative Reynolds number
- ▶ Stokes' viscous drag
- ▶ presence of gravity and diffusivity



# INERTIAL PARTICLES: MODEL AND LIMITATIONS

## Model:

- ▶ given incompressible surrounding flow
- ▶ single spherical particle
- ▶ size smaller than (smallest active) scale of flow
- ▶ low relative Reynolds number
- ▶ Stokes' viscous drag
- ▶ presence of gravity and diffusivity

## Limitations:

- ▶ no interaction with boundaries or other particles

# INERTIAL PARTICLES: MODEL AND LIMITATIONS

## Model:

- ▶ given incompressible surrounding flow
- ▶ single spherical particle
- ▶ size smaller than (smallest active) scale of flow
- ▶ low relative Reynolds number
- ▶ Stokes' viscous drag
- ▶ presence of gravity and diffusivity

## Limitations:

- ▶ no interaction with boundaries or other particles
- ▶ neglect of Basset, Faxen, Oseen, Saffman correction terms

# INERTIAL PARTICLES: MODEL AND LIMITATIONS

## Model:

- ▶ given incompressible surrounding flow
- ▶ single spherical particle
- ▶ size smaller than (smallest active) scale of flow
- ▶ low relative Reynolds number
- ▶ Stokes' viscous drag
- ▶ presence of gravity and diffusivity

## Limitations:

- ▶ no interaction with boundaries or other particles
- ▶ neglect of Basset, Faxen, Oseen, Saffman correction terms
- ▶ effects of non-sphericity, rotationality and high velocity?

# INERTIAL PARTICLES: EQUATIONS

# INERTIAL PARTICLES: EQUATIONS

$$\begin{cases} \dot{\mathbf{x}} = \mathbf{v} + \beta \mathbf{u}[\mathbf{x}(t), t] \\ \dot{\mathbf{v}} = -\frac{\mathbf{v} - (1 - \beta)\mathbf{u}[\mathbf{x}(t), t]}{\tau} + \frac{\sqrt{2\kappa}}{\tau} \boldsymbol{\eta}(t) + (1 - \beta)\mathbf{g} \end{cases}$$

# INERTIAL PARTICLES: EQUATIONS

$$\begin{cases} \dot{\mathbf{x}} = \mathbf{v} + \beta \mathbf{u}[\mathbf{x}(t), t] \\ \dot{\mathbf{v}} = -\frac{\mathbf{v} - (1 - \beta)\mathbf{u}[\mathbf{x}(t), t]}{\tau} + \frac{\sqrt{2\kappa}}{\tau} \boldsymbol{\eta}(t) + (1 - \beta)\mathbf{g} \end{cases}$$

Adimensional numbers:

# INERTIAL PARTICLES: EQUATIONS

$$\begin{cases} \dot{\mathbf{x}} = \mathbf{v} + \beta \mathbf{u}[\mathbf{x}(t), t] \\ \dot{\mathbf{v}} = -\frac{\mathbf{v} - (1 - \beta)\mathbf{u}[\mathbf{x}(t), t]}{\tau} + \frac{\sqrt{2\kappa}}{\tau}\boldsymbol{\eta}(t) + (1 - \beta)\mathbf{g} \end{cases}$$

Adimensional numbers:

$$\blacktriangleright \beta \equiv \frac{3\rho_f}{\rho_f + 2\rho_p}$$

# INERTIAL PARTICLES: EQUATIONS

$$\begin{cases} \dot{\mathbf{x}} = \mathbf{v} + \beta \mathbf{u}[\mathbf{x}(t), t] \\ \dot{\mathbf{v}} = -\frac{\mathbf{v} - (1 - \beta)\mathbf{u}[\mathbf{x}(t), t]}{\tau} + \frac{\sqrt{2\kappa}}{\tau} \boldsymbol{\eta}(t) + (1 - \beta)\mathbf{g} \end{cases}$$

Adimensional numbers:

► $\beta \equiv \frac{3\rho_f}{\rho_f + 2\rho_p}$	$\beta$	0
	particles:	heavy



# INERTIAL PARTICLES: EQUATIONS

$$\begin{cases} \dot{\mathbf{x}} = \mathbf{v} + \beta \mathbf{u}[\mathbf{x}(t), t] \\ \dot{\mathbf{v}} = -\frac{\mathbf{v} - (1 - \beta)\mathbf{u}[\mathbf{x}(t), t]}{\tau} + \frac{\sqrt{2\kappa}}{\tau} \boldsymbol{\eta}(t) + (1 - \beta)\mathbf{g} \end{cases}$$

Adimensional numbers:

►  $\beta \equiv \frac{3\rho_f}{\rho_f + 2\rho_p}$

$\beta$	0	3
particles:	heavy	light

# INERTIAL PARTICLES: EQUATIONS

$$\begin{cases} \dot{\mathbf{x}} = \mathbf{v} + \beta \mathbf{u}[\mathbf{x}(t), t] \\ \dot{\mathbf{v}} = -\frac{\mathbf{v} - (1 - \beta)\mathbf{u}[\mathbf{x}(t), t]}{\tau} + \frac{\sqrt{2\kappa}}{\tau} \boldsymbol{\eta}(t) + (1 - \beta)\mathbf{g} \end{cases}$$

Adimensional numbers:

►  $\beta \equiv \frac{3\rho_f}{\rho_f + 2\rho_p}$

$\beta$	0	$\longleftrightarrow$	1	$\longleftrightarrow$	3
particles:	heavy		tracer		light

# INERTIAL PARTICLES: EQUATIONS

$$\begin{cases} \dot{\mathbf{x}} = \mathbf{v} + \beta \mathbf{u}[\mathbf{x}(t), t] \\ \dot{\mathbf{v}} = -\frac{\mathbf{v} - (1 - \beta)\mathbf{u}[\mathbf{x}(t), t]}{\tau} + \frac{\sqrt{2\kappa}}{\tau} \boldsymbol{\eta}(t) + (1 - \beta)\mathbf{g} \end{cases}$$

Adimensional numbers:

►  $\beta \equiv \frac{3\rho_f}{\rho_f + 2\rho_p}$

$\beta$	0	$\longleftrightarrow$	1	$\longleftrightarrow$	3
particles:	heavy		tracer		light

► **Stokes:**  $St = \frac{\tau}{L/U}$  (inertia)

# INERTIAL PARTICLES: EQUATIONS

$$\begin{cases} \dot{\mathbf{x}} = \mathbf{v} + \beta \mathbf{u}[\mathbf{x}(t), t] \\ \dot{\mathbf{v}} = -\frac{\mathbf{v} - (1 - \beta)\mathbf{u}[\mathbf{x}(t), t]}{\tau} + \frac{\sqrt{2\kappa}}{\tau} \boldsymbol{\eta}(t) + (1 - \beta)\mathbf{g} \end{cases}$$

Adimensional numbers:

►  $\beta \equiv \frac{3\rho_f}{\rho_f + 2\rho_p}$

$\beta$	0	$\longleftrightarrow$	1	$\longleftrightarrow$	3
particles:	heavy		tracer		light

► Stokes:  $St = \frac{\tau}{L/U}$  (inertia)

► **Péclet:**  $Pe = \frac{LU}{\kappa}$  (diffusivity)<sup>-1</sup>

# INERTIAL PARTICLES: EQUATIONS

$$\begin{cases} \dot{\mathbf{x}} = \mathbf{v} + \beta \mathbf{u}[\mathbf{x}(t), t] \\ \dot{\mathbf{v}} = -\frac{\mathbf{v} - (1 - \beta)\mathbf{u}[\mathbf{x}(t), t]}{\tau} + \frac{\sqrt{2\kappa}}{\tau} \boldsymbol{\eta}(t) + (1 - \beta)\mathbf{g} \end{cases}$$

Adimensional numbers:

►  $\beta \equiv \frac{3\rho_f}{\rho_f + 2\rho_p}$

$\beta$	0	$\longleftrightarrow$	1	$\longleftrightarrow$	3
particles:	heavy		tracer		light

► Stokes:  $St = \frac{\tau}{L/U}$  (inertia)

► Péclet:  $Pe = \frac{LU}{\kappa}$  (diffusivity)<sup>-1</sup>

► Froude:  $Fr = \frac{U}{\sqrt{gL}}$  (gravity)<sup>-1/2</sup>

# CONTINUUM VS. DYNAMICAL APPROACH

# CONTINUUM VS. DYNAMICAL APPROACH

- ▶ Continuum description:

# CONTINUUM VS. DYNAMICAL APPROACH

- ▶ Continuum description:

$$\mathbf{v}(\mathbf{x}, t) = (1 - \beta)\mathbf{u}(\mathbf{x}, t) + \tau(1 - \beta)[\mathbf{g} - (\partial_t \mathbf{u} + \mathbf{u} \cdot \partial \mathbf{u})] + O(\tau^2)$$



# CONTINUUM VS. DYNAMICAL APPROACH

- ▶ Continuum description:

$$\mathbf{v}(\mathbf{x}, t) = (1 - \beta)\mathbf{u}(\mathbf{x}, t) + \tau(1 - \beta)[\mathbf{g} - (\partial_t \mathbf{u} + \mathbf{u} \cdot \partial \mathbf{u})] + O(\tau^2)$$

- ▶ Fokker–Planck equation for phase-space density  $\rho(\mathbf{x}, \mathbf{v}, t)$ :

# CONTINUUM VS. DYNAMICAL APPROACH

- ▶ Continuum description:

$$\mathbf{v}(\mathbf{x}, t) = (1 - \beta)\mathbf{u}(\mathbf{x}, t) + \tau(1 - \beta)[\mathbf{g} - (\partial_t \mathbf{u} + \mathbf{u} \cdot \partial \mathbf{u})] + O(\tau^2)$$

- ▶ Fokker–Planck equation for phase-space density  $\rho(\mathbf{x}, \mathbf{v}, t)$ :

$$\frac{\partial}{\partial t} \rho = \frac{\partial}{\partial x_i} (\dots \rho) + \frac{\partial}{\partial v_j} (\dots \rho) + \frac{\kappa}{\tau^2} \frac{\partial^2}{\partial v_k \partial v_k} \rho$$

# CONTINUUM VS. DYNAMICAL APPROACH

- ▶ Continuum description:

$$\mathbf{v}(\mathbf{x}, t) = (1 - \beta)\mathbf{u}(\mathbf{x}, t) + \tau(1 - \beta)[\mathbf{g} - (\partial_t \mathbf{u} + \mathbf{u} \cdot \partial \mathbf{u})] + O(\tau^2)$$

- ▶ Fokker–Planck equation for phase-space density  $\rho(\mathbf{x}, \mathbf{v}, t)$ :

$$\begin{aligned}\frac{\partial}{\partial t} \rho &= \frac{\partial}{\partial x_i} (\dots \rho) + \frac{\partial}{\partial v_j} (\dots \rho) + \frac{\kappa}{\tau^2} \frac{\partial^2}{\partial v_k \partial v_k} \rho \\ &\equiv -\mathcal{L} \rho\end{aligned}$$

# CONTINUUM VS. DYNAMICAL APPROACH

- ▶ Continuum description:

$$\mathbf{v}(\mathbf{x}, t) = (1 - \beta)\mathbf{u}(\mathbf{x}, t) + \tau(1 - \beta)[\mathbf{g} - (\partial_t \mathbf{u} + \mathbf{u} \cdot \partial \mathbf{u})] + O(\tau^2)$$

- ▶ Fokker–Planck equation for phase-space density  $\rho(\mathbf{x}, \mathbf{v}, t)$ :

$$\begin{aligned} \frac{\partial}{\partial t} \rho &= \frac{\partial}{\partial x_i} (\dots \rho) + \frac{\partial}{\partial v_j} (\dots \rho) + \frac{\kappa}{\tau^2} \frac{\partial^2}{\partial v_k \partial v_k} \rho \\ &\equiv -\mathcal{L} \rho \end{aligned}$$

- ▶ Multiscale expansion:

slow variables  $\mathbf{X} \equiv \epsilon \mathbf{x}$  &  $T \equiv \epsilon^2 t$  independent  
(covelocity exclusively fast)

# CONTINUUM VS. DYNAMICAL APPROACH

- Continuum description:

$$\mathbf{v}(\mathbf{x}, t) = (1 - \beta)\mathbf{u}(\mathbf{x}, t) + \tau(1 - \beta)[\mathbf{g} - (\partial_t \mathbf{u} + \mathbf{u} \cdot \partial \mathbf{u})] + O(\tau^2)$$

- Fokker–Planck equation for phase-space density  $\rho(\mathbf{x}, \mathbf{v}, t)$ :

$$\begin{aligned}\frac{\partial}{\partial t} \rho &= \frac{\partial}{\partial x_i} (\dots \rho) + \frac{\partial}{\partial v_j} (\dots \rho) + \frac{\kappa}{\tau^2} \frac{\partial^2}{\partial v_k \partial v_k} \rho \\ &\equiv -\mathcal{L} \rho\end{aligned}$$

- Multiscale expansion:

slow variables  $\mathbf{X} \equiv \epsilon \mathbf{x}$  &  $T \equiv \epsilon^2 t$  independent  
(covelocity exclusively fast)

$$\rho(\mathbf{x}, \mathbf{X}, \mathbf{v}, t, T) = \sum_{l=0}^{\infty} \epsilon^l \rho_l$$

# MULTISCALE EXPANSION

# MULTISCALE EXPANSION

- ▶ Order  $\epsilon^0$ :

# MULTISCALE EXPANSION

- Order  $\epsilon^0$ :

$$(\partial_t + \mathcal{L})\rho_0 = 0$$



# MULTISCALE EXPANSION

► Order  $\epsilon^0$ :

$$(\partial_t + \mathcal{L})\rho_0 = 0 \implies \rho_0 = p(\mathbf{x}, \mathbf{v}, t)P(\mathbf{X}, T)$$

# MULTISCALE EXPANSION

- Order  $\epsilon^0$ :

$$(\partial_t + \mathcal{L})\rho_0 = 0 \implies \rho_0 = p(\mathbf{x}, \mathbf{v}, t)P(\mathbf{X}, T) \implies (\partial_t + \mathcal{L})p = 0$$

# MULTISCALE EXPANSION

- ▶ Order  $\epsilon^0$ :

$$(\partial_t + \mathcal{L})\rho_0 = 0 \implies \rho_0 = p(\mathbf{x}, \mathbf{v}, t)P(\mathbf{X}, T) \implies (\partial_t + \mathcal{L})p = 0$$

- ▶ Order  $\epsilon^1$ :

# MULTISCALE EXPANSION

- Order  $\epsilon^0$ :

$$(\partial_t + \mathcal{L})\rho_0 = 0 \implies \rho_0 = p(\mathbf{x}, \mathbf{v}, t)P(\mathbf{X}, T) \implies (\partial_t + \mathcal{L})p = 0$$

- Order  $\epsilon^1$ :

$$\text{terminal velocity } \mathbf{w} = (1 - \beta)\mathbf{g}\tau + \langle \mathbf{u}(\mathbf{x}, t) \rangle_{p(\mathbf{x}, \mathbf{v}, t)}$$

# MULTISCALE EXPANSION

- Order  $\epsilon^0$ :

$$(\partial_t + \mathcal{L})\rho_0 = 0 \implies \rho_0 = p(\mathbf{x}, \mathbf{v}, t)P(\mathbf{X}, T) \implies (\partial_t + \mathcal{L})p = 0$$

- Order  $\epsilon^1$ :

$$\text{terminal velocity } \mathbf{w} = (1 - \beta)\mathbf{g}\tau + \langle \mathbf{u}(\mathbf{x}, t) \rangle_{p(\mathbf{x}, \mathbf{v}, t)}$$

$$\mathbf{w}_* = (1 - \beta)\mathbf{g}\tau$$

# MULTISCALE EXPANSION

- Order  $\epsilon^0$ :

$$(\partial_t + \mathcal{L})\rho_0 = 0 \implies \rho_0 = p(\mathbf{x}, \mathbf{v}, t)P(\mathbf{X}, T) \implies (\partial_t + \mathcal{L})p = 0$$

- Order  $\epsilon^1$ :

$$\text{terminal velocity } \mathbf{w} = (1 - \beta)\mathbf{g}\tau + \langle \mathbf{u}(\mathbf{x}, t) \rangle_{p(\mathbf{x}, \mathbf{v}, t)}$$

$$\mathbf{w}_* = (1 - \beta)\mathbf{g}\tau \implies \Delta \mathbf{w} = \int d\mathbf{x} d\mathbf{v} dt \mathbf{u}(\mathbf{x}, t) p(\mathbf{x}, \mathbf{v}, t)$$

# MULTISCALE EXPANSION

- Order  $\epsilon^0$ :

$$(\partial_t + \mathcal{L})\rho_0 = 0 \implies \rho_0 = p(\mathbf{x}, \mathbf{v}, t)P(\mathbf{X}, T) \implies (\partial_t + \mathcal{L})p = 0$$

- Order  $\epsilon^1$ :

$$\text{terminal velocity } \mathbf{w} = (1 - \beta)\mathbf{g}\tau + \langle \mathbf{u}(\mathbf{x}, t) \rangle_{p(\mathbf{x}, \mathbf{v}, t)}$$

$$\mathbf{w}_* = (1 - \beta)\mathbf{g}\tau \implies \Delta \mathbf{w} = \int d\mathbf{x} d\mathbf{v} dt \mathbf{u}(\mathbf{x}, t) p(\mathbf{x}, \mathbf{v}, t)$$

$$(\text{solvability condition} \Rightarrow \text{FoR with } \mathbf{w} = 0)$$

- Order  $\epsilon^2$ :

$$\frac{\partial}{\partial T} P = K_{ij} \frac{\partial^2}{\partial X_i \partial X_j} P$$

$$\begin{cases} K_{ij} = \langle v_i + \beta u_i(\mathbf{x}, t) \rangle_{q_j(\mathbf{x}, \mathbf{v}, t)} & (\text{effective diffusivity}) \\ (\partial_t + \mathcal{L})\mathbf{q} = [\mathbf{v} + \beta \mathbf{u}(\mathbf{x}, t)]p(\mathbf{x}, \mathbf{v}, t) & (\text{auxiliary equation}) \end{cases}$$

# MULTISCALE EXPANSION

- Order  $\epsilon^0$ :

$$(\partial_t + \mathcal{L})\rho_0 = 0 \implies \rho_0 = p(\mathbf{x}, \mathbf{v}, t)P(\mathbf{X}, T) \implies (\partial_t + \mathcal{L})p = 0$$

- Order  $\epsilon^1$ :

$$\text{terminal velocity } \mathbf{w} = (1 - \beta)\mathbf{g}\tau + \langle \mathbf{u}(\mathbf{x}, t) \rangle_{p(\mathbf{x}, \mathbf{v}, t)}$$

$$\mathbf{w}_* = (1 - \beta)\mathbf{g}\tau \implies \Delta \mathbf{w} = \int d\mathbf{x} d\mathbf{v} dt \mathbf{u}(\mathbf{x}, t) p(\mathbf{x}, \mathbf{v}, t)$$

(solvability condition  $\Rightarrow$  FoR with  $\mathbf{w} = 0$ )

- Order  $\epsilon^2$ :

$$\frac{\partial}{\partial T} P = K_{ij} \frac{\partial^2}{\partial X_i \partial X_j} P$$

$$\begin{cases} K_{ij} = \langle v_i + \beta u_i(\mathbf{x}, t) \rangle_{q_j(\mathbf{x}, \mathbf{v}, t)} & \text{(effective diffusivity)} \\ (\partial_t + \mathcal{L})\mathbf{q} = [\mathbf{v} + \beta \mathbf{u}(\mathbf{x}, t)]p(\mathbf{x}, \mathbf{v}, t) & \text{(auxiliary equation)} \end{cases}$$



# MULTISCALE EXPANSION

- Order  $\epsilon^0$ :

$$(\partial_t + \mathcal{L})\rho_0 = 0 \implies \rho_0 = p(\mathbf{x}, \mathbf{v}, t)P(\mathbf{X}, T) \implies (\partial_t + \mathcal{L})p = 0$$

- Order  $\epsilon^1$ :

$$\text{terminal velocity } \mathbf{w} = (1 - \beta)\mathbf{g}\tau + \langle \mathbf{u}(\mathbf{x}, t) \rangle_{p(\mathbf{x}, \mathbf{v}, t)}$$

$$\mathbf{w}_* = (1 - \beta)\mathbf{g}\tau \implies \Delta \mathbf{w} = \int d\mathbf{x} d\mathbf{v} dt \mathbf{u}(\mathbf{x}, t) p(\mathbf{x}, \mathbf{v}, t)$$

(solvability condition  $\Rightarrow$  FoR with  $\mathbf{w} = 0$ )

- Order  $\epsilon^2$ :

$$\frac{\partial}{\partial T} P = K_{ij} \frac{\partial^2}{\partial X_i \partial X_j} P$$

$$\begin{cases} K_{ij} = \langle v_i + \beta u_i(\mathbf{x}, t) \rangle_{q_j(\mathbf{x}, \mathbf{v}, t)} & \text{(effective diffusivity)} \\ (\partial_t + \mathcal{L})\mathbf{q} = [\mathbf{v} + \beta \mathbf{u}(\mathbf{x}, t)]p(\mathbf{x}, \mathbf{v}, t) & \text{(auxiliary equation)} \end{cases}$$

# SECOND-QUANTIZATION FORMALISM

# SECOND-QUANTIZATION FORMALISM

- ▶ adimensional reformulation of  $(\partial_t + \mathcal{L})p = 0$ :  
operator with terms  $S t^{-1}$ ,  $S t^{-1/2}$ ,  $S t^0$ ,  $S t^{1/2}$

# SECOND-QUANTIZATION FORMALISM

- ▶ adimensional reformulation of  $(\partial_t + \mathcal{L})p = 0$ :  
operator with terms  $St^{-1}$ ,  $St^{-1/2}$ ,  $St^0$ ,  $St^{1/2}$
- ▶ expansion  $p(\mathbf{x}, \mathbf{v}, t) = \sum_{n=0}^{\infty} St^{n/2} p_n(\mathbf{x}, \mathbf{v}, t)$

# SECOND-QUANTIZATION FORMALISM

- ▶ adimensional reformulation of  $(\partial_t + \mathcal{L})p = 0$ :  
operator with terms  $St^{-1}$ ,  $St^{-1/2}$ ,  $St^0$ ,  $St^{1/2}$
- ▶ expansion  $p(\mathbf{x}, \mathbf{v}, t) = \sum_{n=0}^{\infty} St^{n/2} p_n(\mathbf{x}, \mathbf{v}, t)$
- ▶ Hermitian reformulation and introduction of vacuum, creation, annihilation for  $\mathbf{v}$

# SECOND-QUANTIZATION FORMALISM

- ▶ adimensional reformulation of  $(\partial_t + \mathcal{L})p = 0$ :  
operator with terms  $St^{-1}$ ,  $St^{-1/2}$ ,  $St^0$ ,  $St^{1/2}$
- ▶ expansion  $p(\mathbf{x}, \mathbf{v}, t) = \sum_{n=0}^{\infty} St^{n/2} p_n(\mathbf{x}, \mathbf{v}, t)$
- ▶ Hermitian reformulation and introduction of vacuum, creation, annihilation for  $\mathbf{v}$
- ▶  $(\partial_t + \mathbf{u} \cdot \partial - \kappa \partial^2) p_n = f(p_{n-2}, p_{n-4}, \dots)$  with  $p_0 = 1$ ,  $p_1 = 0$

# SECOND-QUANTIZATION FORMALISM

- ▶ adimensional reformulation of  $(\partial_t + \mathcal{L})p = 0$ :  
operator with terms  $St^{-1}$ ,  $St^{-1/2}$ ,  $St^0$ ,  $St^{1/2}$
- ▶ expansion  $p(\mathbf{x}, \mathbf{v}, t) = \sum_{n=0}^{\infty} St^{n/2} p_n(\mathbf{x}, \mathbf{v}, t)$
- ▶ Hermitian reformulation and introduction of vacuum, creation, annihilation for  $\mathbf{v}$
- ▶  $(\partial_t + \mathbf{u} \cdot \partial - \kappa \partial^2) p_n = f(p_{n-2}, p_{n-4}, \dots)$  with  $p_0 = 1$ ,  $p_1 = 0$
- ▶ bare terminal velocity  $w_* = (1 - \beta)Fr^{-2}St$

# SECOND-QUANTIZATION FORMALISM

- ▶ adimensional reformulation of  $(\partial_t + \mathcal{L})p = 0$ :  
operator with terms  $\text{St}^{-1}$ ,  $\text{St}^{-1/2}$ ,  $\text{St}^0$ ,  $\text{St}^{1/2}$
- ▶ expansion  $p(\mathbf{x}, \mathbf{v}, t) = \sum_{n=0}^{\infty} \text{St}^{n/2} p_n(\mathbf{x}, \mathbf{v}, t)$
- ▶ Hermitian reformulation and introduction of vacuum, creation, annihilation for  $\mathbf{v}$
- ▶  $(\partial_t + \mathbf{u} \cdot \partial - \kappa \partial^2) p_n = f(p_{n-2}, p_{n-4}, \dots)$  with  $p_0 = 1$ ,  $p_1 = 0$
- ▶ bare terminal velocity  $w_* = (1 - \beta) \text{Fr}^{-2} \text{St}$   
 $\implies \Delta w = \sum_{m=0}^{\infty} \text{St}^{2+m} \int d\mathbf{x} dt \mathbf{u}(\mathbf{x}, t) p_{4+2m}(\mathbf{x}, t)$



# SECOND-QUANTIZATION FORMALISM

- ▶ adimensional reformulation of  $(\partial_t + \mathcal{L})p = 0$ :  
operator with terms  $St^{-1}$ ,  $St^{-1/2}$ ,  $St^0$ ,  $St^{1/2}$
- ▶ expansion  $p(\mathbf{x}, \mathbf{v}, t) = \sum_{n=0}^{\infty} St^{n/2} p_n(\mathbf{x}, \mathbf{v}, t)$
- ▶ Hermitian reformulation and introduction of vacuum, creation, annihilation for  $\mathbf{v}$
- ▶  $(\partial_t + \mathbf{u} \cdot \partial - \kappa \partial^2) p_n = f(p_{n-2}, p_{n-4}, \dots)$  with  $p_0 = 1$ ,  $p_1 = 0$
- ▶ bare terminal velocity  $w_* = (1 - \beta) Fr^{-2} St$   
 $\implies \Delta w = \sum_{m=0}^{\infty} St^{2+m} \int d\mathbf{x} dt \mathbf{u}(\mathbf{x}, t) p_{4+2m}(\mathbf{x}, t)$   
 $\implies$  possibility of variation  $O(St^2)$  from bare value

# CORRECTION AT ORDER $St^2$

## CORRECTION AT ORDER $St^2$

- ▶ proportional to gravity ( $\propto Fr^{-2}$ )

## CORRECTION AT ORDER $St^2$

- ▶ proportional to gravity ( $\propto Fr^{-2}$ )
- ▶ same sign for both heavy and light particles ( $\propto (1 - \beta)^2$ )

## CORRECTION AT ORDER $St^2$

- ▶ proportional to gravity ( $\propto Fr^{-2}$ )
- ▶ same sign for both heavy and light particles ( $\propto (1 - \beta)^2$ )  
 $\implies$  acceleration in falling and slowdown in rising, or viceversa

## CORRECTION AT ORDER $St^2$

- ▶ proportional to gravity ( $\propto Fr^{-2}$ )
- ▶ same sign for both heavy and light particles ( $\propto (1 - \beta)^2$ )  
 $\implies$  acceleration in falling and slowdown in rising, or viceversa
- ▶ actual presence and sign: dependent on the flow

## CORRECTION AT ORDER $St^2$

- ▶ proportional to gravity ( $\propto Fr^{-2}$ )
- ▶ same sign for both heavy and light particles ( $\propto (1 - \beta)^2$ )  
 $\implies$  acceleration in falling and slowdown in rising, or viceversa
- ▶ actual presence and sign: dependent on the flow

Numerical examples:

# CORRECTION AT ORDER $St^2$

- ▶ proportional to gravity ( $\propto Fr^{-2}$ )
- ▶ same sign for both heavy and light particles ( $\propto (1 - \beta)^2$ )  
 $\implies$  **acceleration in falling** and slowdown in rising, or viceversa
- ▶ actual presence and sign: dependent on the flow

Numerical examples:

- ▶ **ABC (3D) or BC (2D) flows**



# CORRECTION AT ORDER $St^2$

- ▶ proportional to gravity ( $\propto Fr^{-2}$ )
- ▶ same sign for both heavy and light particles ( $\propto (1 - \beta)^2$ )  
 $\implies$  acceleration in falling and slowdown in rising, or viceversa
- ▶ actual presence and sign: dependent on the flow

Numerical examples:

- ▶ ABC (3D) or BC (2D) flows
- ▶ Gollub flow (2D) with  $\mathbf{g} \parallel -x_2$   
$$\mathbf{u} = (\sin(kx_1) \cos[x_2 + \sin(\omega t)], -\cos(kx_1) \sin[x_2 + \sin(\omega t)])$$

# CORRECTION AT ORDER $St^2$

- ▶ proportional to gravity ( $\propto Fr^{-2}$ )
- ▶ same sign for both heavy and light particles ( $\propto (1 - \beta)^2$ )  
 $\implies$  acceleration in falling and slowdown in rising, or viceversa
- ▶ actual presence and sign: dependent on the flow

Numerical examples:

- ▶ ABC (3D) or BC (2D) flows
- ▶ Gollub flow (2D) with  $\mathbf{g} \parallel -x_2$   
$$\mathbf{u} = (\sin(kx_1) \cos[x_2 + \sin(\omega t)], -\cos(kx_1) \sin[x_2 + \sin(\omega t)])$$
  
 $\implies$  dependence on oscillation and aspect ratio:

# CORRECTION AT ORDER $St^2$

- ▶ proportional to gravity ( $\propto Fr^{-2}$ )
- ▶ same sign for both heavy and light particles ( $\propto (1 - \beta)^2$ )  
 $\implies$  **acceleration in falling** and slowdown in rising, or viceversa
- ▶ actual presence and sign: dependent on the flow

Numerical examples:

- ▶ ABC (3D) or BC (2D) flows
- ▶ Gollub flow (2D) with  $\mathbf{g} \parallel -x_2$   
$$\mathbf{u} = (\sin(kx_1) \cos[x_2 + \sin(\omega t)], -\cos(kx_1) \sin[x_2 + \sin(\omega t)])$$
  
 $\implies$  dependence on oscillation and aspect ratio:

$$\omega = 0, k = 1$$

# CORRECTION AT ORDER $St^2$

- ▶ proportional to gravity ( $\propto Fr^{-2}$ )
- ▶ same sign for both heavy and light particles ( $\propto (1 - \beta)^2$ )  
 $\implies$  acceleration in falling and slowdown in rising, or **viceversa**
- ▶ actual presence and sign: dependent on the flow

Numerical examples:

- ▶ ABC (3D) or BC (2D) flows
- ▶ Gollub flow (2D) with  $\mathbf{g} \parallel -x_2$   
 $\mathbf{u} = (\sin(kx_1) \cos[x_2 + \sin(\omega t)], -\cos(kx_1) \sin[x_2 + \sin(\omega t)])$   
 $\implies$  dependence on oscillation and aspect ratio:

$$\omega = 0, k = 1 \longleftrightarrow \left\{ \begin{array}{l} \omega \neq 0, k = 1 \end{array} \right.$$

# CORRECTION AT ORDER $St^2$

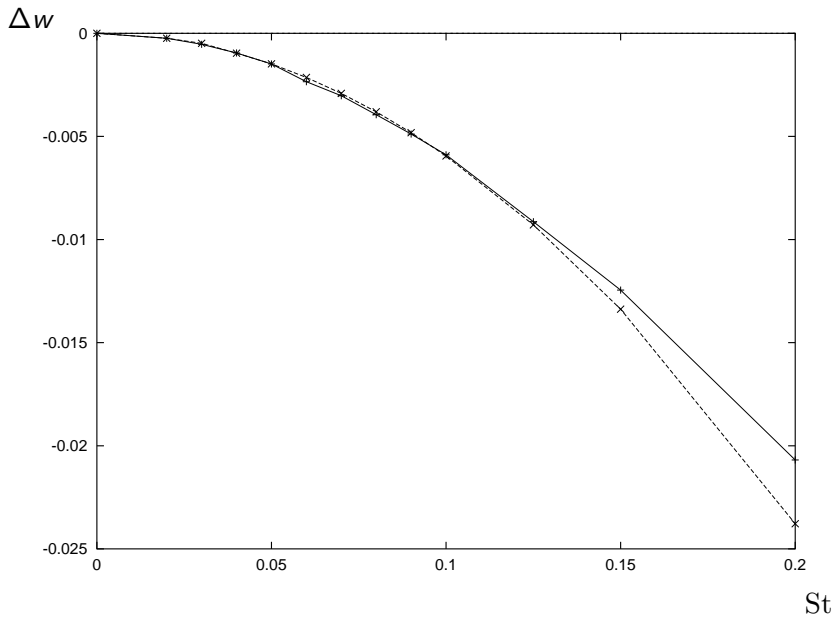
- ▶ proportional to gravity ( $\propto Fr^{-2}$ )
- ▶ same sign for both heavy and light particles ( $\propto (1 - \beta)^2$ )  
 $\implies$  acceleration in falling and slowdown in rising, or **viceversa**
- ▶ actual presence and sign: dependent on the flow

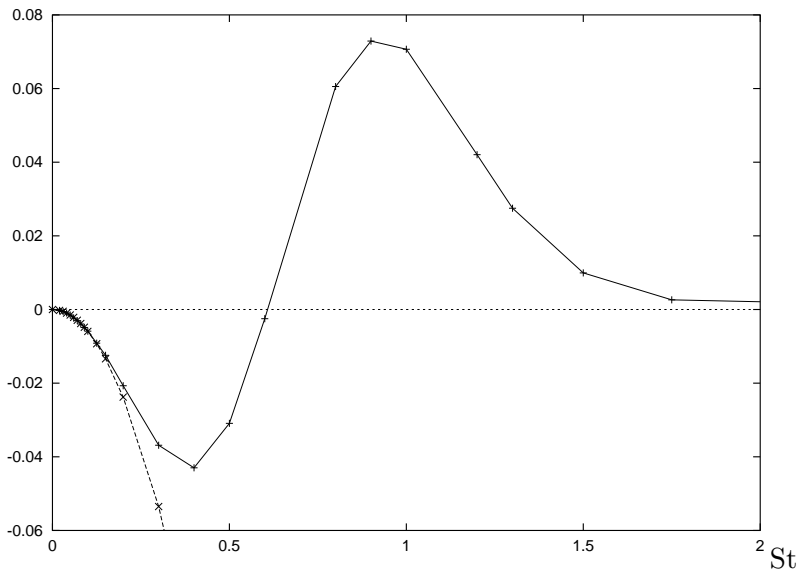
Numerical examples:

- ▶ ABC (3D) or BC (2D) flows
- ▶ Gollub flow (2D) with  $\mathbf{g} \parallel -x_2$   
 $\mathbf{u} = (\sin(kx_1) \cos[x_2 + \sin(\omega t)], -\cos(kx_1) \sin[x_2 + \sin(\omega t)])$

$\implies$  dependence on oscillation and aspect ratio:

$$\omega = 0, k = 1 \longleftrightarrow \begin{cases} \omega \neq 0, k = 1 \\ \omega = 0, k \neq 1 \end{cases}$$



$\Delta w$ 

# SINGULAR EXPANSIONS AND DIFFUSIVITY



# SINGULAR EXPANSIONS AND DIFFUSIVITY

- Expansion at large  $St$  (ballistic case):

# SINGULAR EXPANSIONS AND DIFFUSIVITY

- Expansion at large  $St$  (ballistic case):  
singular (white  $\rightarrow$  coloured noise)

# SINGULAR EXPANSIONS AND DIFFUSIVITY

- ▶ Expansion at large  $St$  (ballistic case):  
singular (white  $\rightarrow$  coloured noise)
- ▶ Dependence on diffusivity:

# SINGULAR EXPANSIONS AND DIFFUSIVITY

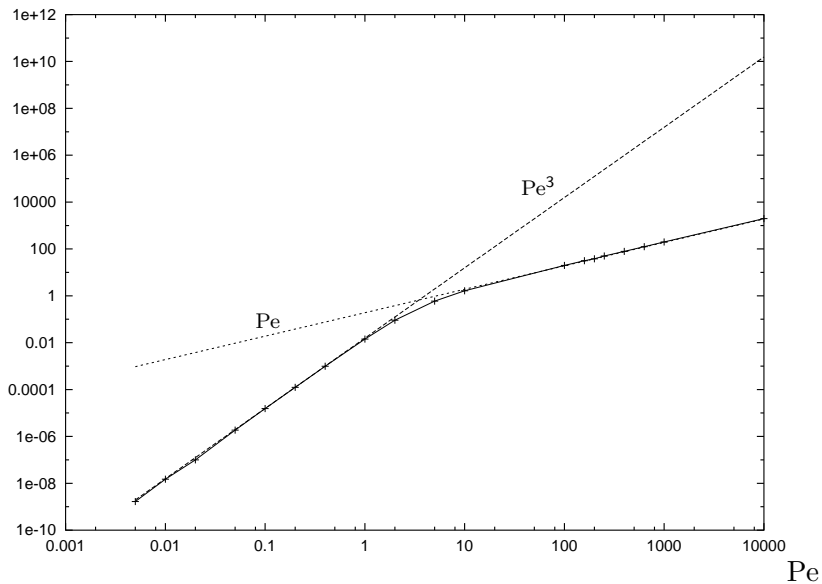
- ▶ Expansion at large  $St$  (ballistic case):  
singular (white  $\rightarrow$  coloured noise)
- ▶ Dependence on diffusivity:
  - ▶ expansion at large  $Pe$ : singular

# SINGULAR EXPANSIONS AND DIFFUSIVITY

- ▶ Expansion at large  $St$  (ballistic case):  
singular (white  $\rightarrow$  coloured noise)
- ▶ Dependence on diffusivity:
  - ▶ expansion at large  $Pe$ : singular
  - ▶ expansion at small  $Pe$ :  $O(Pe^2)$  or higher

# SINGULAR EXPANSIONS AND DIFFUSIVITY

- ▶ Expansion at large  $St$  (ballistic case):  
singular (white  $\rightarrow$  coloured noise)
- ▶ Dependence on diffusivity:
  - ▶ expansion at large  $Pe$ : singular
  - ▶ expansion at small  $Pe$ :  $O(Pe^2)$  or higher
  - ▶ special case (stationary Gollub f.): behaviours  $Pe$  and  $Pe^3$  respectively

$\Delta w$ 

# FURTHER DEVELOPMENTS



# FURTHER DEVELOPMENTS

- ▶ Study at large (and intermediate)  $St$  and large  $Pe$

# FURTHER DEVELOPMENTS

- ▶ Study at large (and intermediate)  $St$  and large  $Pe$
- ▶ (with P. Olla)  
Turbulent, one-dimensional Gaussian flow, at large  $St$   
or small  $Fr$  (absence of  $Pe$  but appearance of  $Ku$ ):

# FURTHER DEVELOPMENTS

- ▶ Study at large (and intermediate)  $St$  and large  $Pe$
- ▶ (with P. Olla)  
Turbulent, one-dimensional Gaussian flow, at large  $St$   
or small  $Fr$  (absence of  $Pe$  but appearance of  $Ku$ ):  
slowdown in falling for heavy p.

# FURTHER DEVELOPMENTS

- ▶ Study at large (and intermediate)  $St$  and large  $Pe$
- ▶ (with P. Olla)  
Turbulent, one-dimensional Gaussian flow, at large  $St$  or small  $Fr$  (absence of  $Pe$  but appearance of  $Ku$ ):  
slowdown in falling for heavy p.
- ▶ Study of effective diffusivity